

論文 / 著書情報  
Article / Book Information

Title	Sub-millimeter propagation of antiferromagnetic magnons via magnon-photon coupling
Authors	Ryo Kainuma, Keita Matsumoto, Toshimitsu Ito, Takuya Satoh
Citation	npj Spintronics, Vol. 2, , pp. 31-1-5
Pub. date	2024, 7
DOI	<a href="https://doi.org/10.1038/s44306-024-00034-3">https://doi.org/10.1038/s44306-024-00034-3</a>
Creative Commons	Information is in the article.
Note	The version of record of this article, first published in npj Spintronics, is available online at Publisher ' s website: <a href="https://doi.org/10.1038/s44306-024-00034-3">https://doi.org/10.1038/s44306-024-00034-3</a>

<https://doi.org/10.1038/s44306-024-00034-3>

# Sub-millimeter propagation of antiferromagnetic magnons via magnon-photon coupling

Check for updates

Ryo Kainuma<sup>1,5</sup>, Keita Matsumoto<sup>1,2,5</sup>, Toshimitsu Ito<sup>3</sup> & Takuya Satoh<sup>1,4</sup>✉

For the realization of magnon-based current-free technologies, referred to as magnonics, all-optical control of magnons is an important technique for both fundamental research and practical applications. Magnon-polariton is a coupled state of magnon and photon in a magnetic medium, expected to exhibit magnon-like controllability and photon-like high-speed propagation. While recent studies have observed magnon-polaritons as modulation of incident terahertz waves, the influence of magnon-photon coupling on magnon propagation properties remains unexplored. This study aimed to observe the spatiotemporal dynamics of coherent magnon-polaritons through time-resolved imaging measurements. BiFeO<sub>3</sub> was selected as the sample due to its anticipated strong coupling between magnons and photons. The observed dynamics suggest that antiferromagnetic magnons can propagate over long distances, up to hundreds of micrometers, through strong coupling with photons. These results enhance our understanding of the optical control of magnonic systems, thereby paving the way for terahertz opto-magnonics.

In the past decade, significant research has been conducted on revolutionary information-processing devices that utilize magnons as an information carrier, marking significant progress in the field of magnonics<sup>1,2</sup>. Following the study reporting the spin-wave logic gate<sup>3</sup>, there has been extensive exploration of ferromagnetic magnons. In recent years, research focus has extended from ferromagnets (FMs) to antiferromagnets (AFMs)<sup>4,5</sup>. This extension is primarily due to the higher resonance frequencies of AFMs, extending into the terahertz region. It is desired that magnons are excited coherently and propagate over distances longer than their wavelengths to leverage superposition and nonlinearity<sup>2,6</sup>. However, the excitation and detection of magnons in AFMs pose greater challenges compared to FMs due to the absence of net magnetization in AFMs.

Established methods for magnon excitation include electrical spin injection via the spin Hall effect<sup>7</sup> and the use of nanometric coplanar waveguides. While these methods can be applied to AFMs<sup>8,9</sup>, they require microfabrication. On the other hand, ultrashort light pulses can excite and detect coherent magnons in a noncontact manner without the need for microfabrication. Previous studies have demonstrated the all-optical excitation and detection of antiferromagnetic magnons, attributing these optical excitation mechanisms to inverse magneto-optical effects<sup>10,11</sup>. Such nonthermal optical excitation is typically limited to addressing magnons near the Brillouin zone center or edge, as magnons

in these regions often have a group velocity close to zero in AFMs, hindering their propagation.

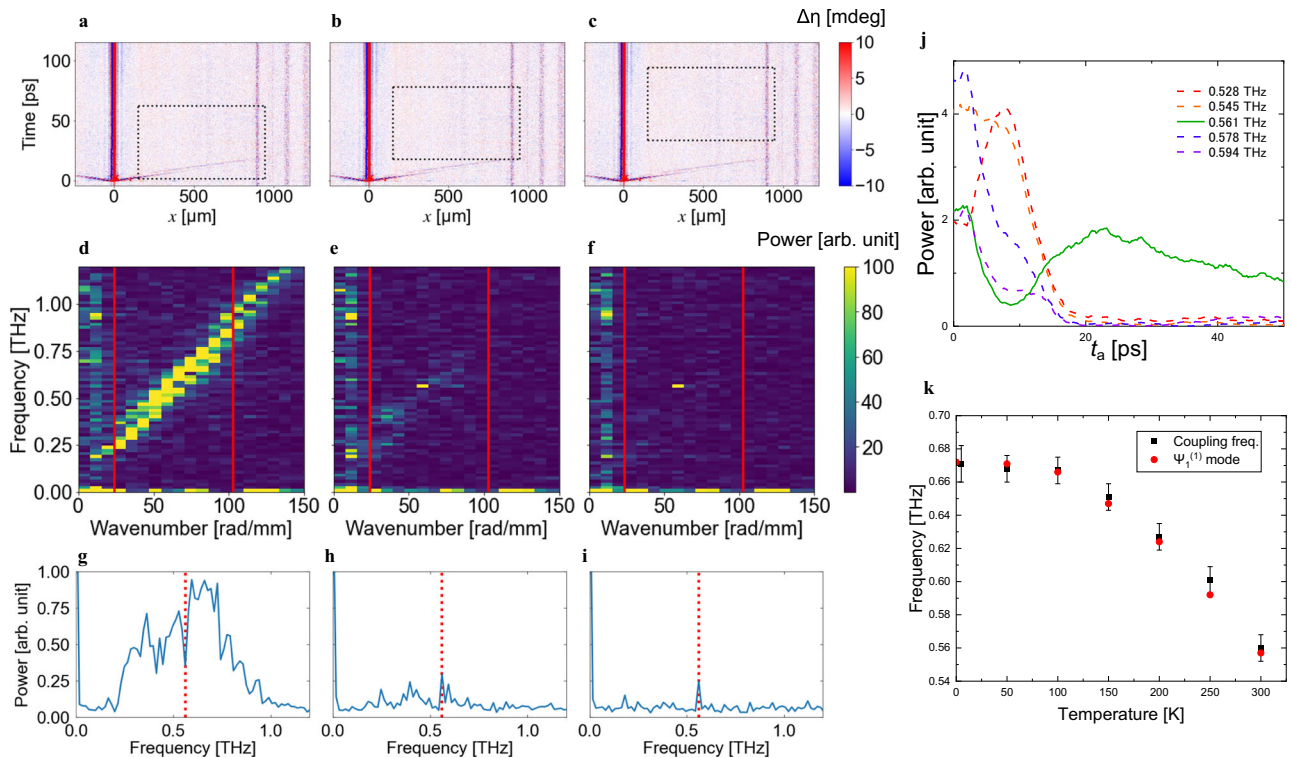
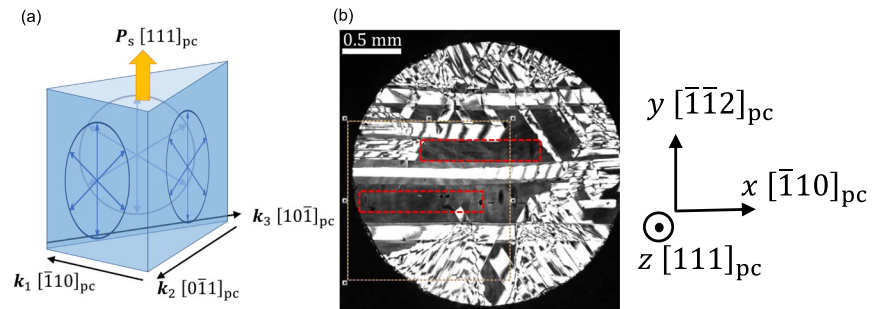
The energy transfer between magnons and photons gives rise to their coupled states, known as magnon-polaritons. Magnon-polaritons can be optically excited and are expected to exhibit a nonzero group velocity due to their anti-cross-energy dispersion. Previous studies have primarily focused on the magnonic features of the transmitted terahertz wave, which lacks spatial resolution<sup>12–14</sup>. Consequently, the influence of magnon-photon coupling on magnon propagation properties remains unclear.

This article reports on the all-optical coherent excitation and detection of magnon-polaritons in BiFeO<sub>3</sub>. This material is anticipated to demonstrate strong magnon-photon coupling due to the presence of electromagnons<sup>15</sup>. Additionally, the multiferroic nature of BiFeO<sub>3</sub> (ref. 16) enabled the excitation and detection of magnon-polaritons in a single material in contrast to a previous study using a hybrid ferroelectric-antiferromagnetic waveguide<sup>17</sup>. Time-resolved imaging measurements revealed that antiferromagnetic magnon in BiFeO<sub>3</sub> can propagate over long distances, up to hundreds of micrometers in tens of picoseconds, by forming magnon-polaritons. The excitation is based on nonthermal mechanism, which does not require optical absorption<sup>18</sup>.

Figure 1a depicts a schematic of the multiferroic order in BiFeO<sub>3</sub>. At room temperature, BiFeO<sub>3</sub> exhibits multiferroicity, characterized by

<sup>1</sup>Department of Physics, Tokyo Institute of Technology, Tokyo 152-8551, Japan. <sup>2</sup>Department of Physics, Kyushu University, Fukuoka 819-0395, Japan. <sup>3</sup>National Institute of Advanced Industrial Science and Technology, Tsukuba 305-8565, Japan. <sup>4</sup>Quantum Research Center for Chirality, Institute for Molecular Science, Okazaki 444-8585, Japan. <sup>5</sup>These authors contributed equally: Ryo Kainuma, Keita Matsumoto. ✉e-mail: [satoh@phys.titech.ac.jp](mailto:satoh@phys.titech.ac.jp)

**Fig. 1 | Multiferroic BiFeO<sub>3</sub> single crystal.** **a** The schematic structure of multiferroic order of BiFeO<sub>3</sub>. The blue bidirectional arrows illustrate the three directional magnetic domains of the cycloidal order. **b** BiFeO<sub>3</sub> single crystal used for our measurements. The spatiotemporal waveforms observed in the two red boxes were utilized for the analysis. The yellow dashed square represents the area of 1000 × 1000 pixels that the camera measures at a time.



**Fig. 2 | Coherent propagation of magnon-polaritons and phonon-polaritons.** **a** The spatiotemporal waveform observed in the pump-probe imaging measurement. The colors in the heat map, ranging from red to blue, represent the ellipticity change of the circularly polarized probe pulse induced at each point of the sample. The black dashed frame delineates the region of the Fourier transform for **(d)**. **b, c** shows the same waveform as **(a)** with different Fourier transform regions for **(e, f)**, respectively. **g–i** are distributions in **(d–f)** integrated in the wavenumber

direction within the region delimited by the red lines in **(d–f)**. The red dashed line in **(g–i)** represents 0.56 THz. **j** represents the time evolution of the power of the components around the (anti-)crossing point. **k** illustrates the temperature dependence of the frequency, which exhibited anomalous behavior in **(j)**, along with the eigenfrequency of the magnon  $\Psi_1^{(1)}$  mode<sup>28</sup>. The error bars are defined by the frequency resolutions.

ferroelectricity ( $T_C = 1100$  K) and antiferromagnetism ( $T_N = 640$  K). In the antiferromagnetic phase, the Néel vector  $\mathbf{L}$  displays a cycloidal order with a period of 62 nm<sup>19,20</sup>. The wavevector  $\mathbf{q}$  of the cycloidal order typically orients in the  $[\bar{1}10]_{pc}$ ,  $[0\bar{1}1]_{pc}$ ,  $[10\bar{1}]_{pc}$  directions, perpendicular to the ferroelectric polarization,  $\mathbf{P}_s \parallel [111]_{pc}$ . Antiferromagnetic magnons, the eigenmodes of  $\mathbf{L}$ , are classified into  $\Phi$  and  $\Psi$  modes, corresponding to in-plane and out-of-plane oscillations, respectively, relative to the plane spanned by  $\mathbf{P}_s$  and  $\mathbf{q}$  (ref. 21). The spatial periodicity of the cycloidal order folds the magnon Brillouin zone, and deviations can be expanded using modulated plane wave solutions  $\exp[-i(\mathbf{k} + n\mathbf{q}) \cdot \mathbf{r}]$ , where  $n$  denotes the  $n$ th magnon Brillouin zone. This allows magnons with wavevectors  $\mathbf{k} + n\mathbf{q}$  to be accessible via optical measurements, such as Raman<sup>15,22,23</sup> or infrared spectroscopy<sup>24–29</sup> and pump-probe measurements<sup>30</sup>. Accordingly, the magnon eigenmodes are labeled as  $\Phi_n$  and  $\Psi_n$ . Additionally, eigenmodes split by higher-order perturbations are classified using subscripts, as in  $\Phi_n^{(m)}$  and  $\Psi_n^{(m)}$  (ref. 31).

## Results

The BiFeO<sub>3</sub> single crystal was cut along the  $(111)_{pc}$  plane. Figure 1b displays the ferroelectric domain structure of the BiFeO<sub>3</sub> sample, observed using a CMOS camera under crossed-Nicols conditions. The gray regions within the red rectangular frames did not exhibit optical birefringence, indicating that the polarization ( $\mathbf{P}_s$ ) was homogeneously perpendicular to the surface in this area. The region inside the frame was selected for the observations.

Figure 2a depicts the spatiotemporal waveform propagating in the  $x$ -direction, obtained by integrating the entire waveform (see Supplementary Movie SM1) along the  $y$ -axis. It is evident that the wavepackets propagated beyond the pump spot at  $x = 0 \mu\text{m}$ . These wavepackets corresponded to phonon-dressed terahertz electromagnetic waves, also known as phonon-polaritons<sup>32</sup>. The excitation mechanism of phonon-polaritons is based on difference-frequency generation<sup>33</sup> (see also Supplementary Information I).

A Fourier transform was performed to visualize the population of phonon-polaritons and magnon-polaritons in  $k$ - $f$  space, employing waveforms ranging from 150 to 945  $\mu\text{m}$  from the excitation spot (the region within the black dashed frame, as depicted in Fig. 2a). The frequency resolution is primarily limited by the time window (50 ps). Figure 2d illustrates the linear dispersion of phonon-polaritons. Through this method, components with a finite group velocity were selectively extracted when excluding the excitation spot from the Fourier transform. Consequently, Fig. 2d does not contain signals of phonons or magnons localized at the pump spot. The velocity of the phonon-polaritons was  $\sim 6.1 \times 10^7 \text{ m s}^{-1}$ , leading to a corresponding refractive index of  $n_{\text{THz}} = 4.9$ . This value aligns with that reported in a previous study [ $n_{\text{THz}} = 5.2$  at 400 K (ref. 34)], and it surpasses the refractive index for the pump pulse ( $n_{\text{pu}} = 2.7$ ). Note that the splitting of the phonon-polariton mode is visible in Fig. 2d, which may be attributed to the multimodal nature of the phonon-polariton with its finite thickness<sup>35</sup>. However, anticrossings corresponding to magnon-photon coupling were not clearly identified.

We translated the Fourier transform domain in the time direction to selectively extract the components with different group velocities, as depicted in Fig. 2a–c. Interestingly, we observed that a specific component at  $f = 0.56 \text{ THz}$  and  $k = 60 \text{ rad mm}^{-1}$  exhibited a significant increase in power later than the other components, as illustrated in Fig. 2d–f in sequence (refer to Supplementary Movie SM2). Subsequent temperature-dependence measurements revealed that the frequencies at which these specific components appeared were consistent with the frequencies of the magnon  $\Psi_1^{(1)}$  mode<sup>25–28,30</sup>, as shown in Fig. 2k (see also Supplementary Information II for the distribution power spectra). Hence, we conclude that this signal is located at the anticrossing point of the dispersion curves of magnon  $\Psi_1^{(1)}$  mode and photons, indicating that the signal originates from magnon-polaritons.

To extract additional information from the heat maps (Fig. 2d–f), we first integrated the excitation intensities near the anticrossing in the  $k$ -direction, resulting in Fig. 2g–i. Figure 2g reveals a dip at 0.56 THz, supporting the notion that the magnon-polariton exhibits a relatively slower group velocity compared to the other nearly non-perturbed photons. Additionally, Fig. 2h, i illustrate that the magnon-polaritons reached the Fourier transform region after the other components had passed through. For further confirmation, we plotted the time evolution of the powers of several frequency components at  $k = 60 \text{ rad mm}^{-1}$ , as shown in Fig. 2j. The horizontal axis represents the start time ( $t_a$ ) of the time domain used in the Fourier transform. The solid and dashed lines represent the time evolution of the power of each frequency component, corresponding to the change in pixel brightness in Fig. 2d–f. As anticipated, a noticeable difference in behavior was observed, with the magnon-polaritons of 0.56 THz component increasing later than those of the other components. These observations confirmed that magnon-polariton exhibit a group velocity in the intermediate region between the non-perturbed magnons and photons (phonon-polaritons).

Magnon-polaritons near the anticrossing point exhibited group velocities in the range of  $0\text{--}6.1 \times 10^7 \text{ m s}^{-1} (= c/n_{\text{THz}})$  depending on the wavenumbers. Assuming an average group velocity of  $3.0 \times 10^7 \text{ m s}^{-1}$  for the excited magnon-polariton, as a value in the middle of this range, it can be concluded that the group velocity of the magnon-polariton is three orders of magnitude higher than that in the conventional ferromagnetic system<sup>36</sup>.

## Discussion

In general, the magnon dispersion of antiferromagnets is described as follows<sup>18</sup>:

$$\omega = \sqrt{\omega_0^2 + (v_0 k)^2} \quad (1)$$

For orthoferrite in ref. 18,  $\omega_0 \approx 2\pi \times 175 \times 10^9 \text{ rad s}^{-1}$  and  $v_0 \approx 2 \times 10^4 \text{ m s}^{-1}$ . The magnon group velocity  $\partial\omega/\partial k$  of  $1.3 \times 10^4 \text{ m s}^{-1}$  was

obtained at a magnon wavenumber of  $k = 4.2 \times 10^4 \text{ rad mm}^{-1}$ . The propagation length can be calculated as the product of the group velocity and the lifetime. With a lifetime of 85 ps, this gives rise to a propagation length of  $\sim 10^{-6} \text{ m}$ . For BiFeO<sub>3</sub> in ref. 22 the magnon group velocity of  $1.4 \times 10^4 \text{ m s}^{-1}$  was obtained at the magnon wavenumber of  $k = q = 2\pi/ (62 \text{ nm}) \approx 10^5 \text{ rad mm}^{-1}$ . In both cases, the group velocity is  $\sim 10^4 \text{ m s}^{-1}$  at a wavenumber of  $\sim 10^5 \text{ rad mm}^{-1}$ . It is important to note that these velocities correspond to pure magnons, not to coupled magnon-polaritons. In our case, the excited wavenumber of the magnon-polariton is  $k \approx 60 \text{ rad mm}^{-1}$ . Without coupling to phonon-polariton, the group velocity of the  $\Psi_1^{(1)}$  magnon at this wavenumber would be  $\sim 10^4 \text{ m s}^{-1}$  for  $k \parallel \mathbf{q}$  and  $\sim 10 \text{ m s}^{-1}$  for  $k \perp \mathbf{q}$ <sup>21</sup>. The lifetime of  $\sim 100 \text{ ps}$  (ref. 30) leads to the propagation length of  $\sim 10^{-6} \text{ m}$  for  $k \parallel \mathbf{q}$  and  $\sim 10^{-9} \text{ m}$  for  $k \perp \mathbf{q}$ . However, when magnons are coupled to phonon-polaritons with a velocity of  $c/n_{\text{THz}} = 6.1 \times 10^7 \text{ m s}^{-1}$ , as observed in our study, the group velocity of the magnon-polariton is significantly enhanced at the anticrossing point, reaching an average of  $3 \times 10^7 \text{ m s}^{-1}$ . The lifetime of 40 ps (as shown in Fig. 2j) gives rise to the propagation length of  $\sim 10^{-3} \text{ m}$ . The sub-millimeter propagation of magnon-polaritons is also evident from Fig. 2d–f, obtained by Fourier transform from 150 to 945  $\mu\text{m}$ , particularly with the presence of the spot at  $k \approx 60 \text{ rad mm}^{-1}$  associated with magnon-polaritons in Fig. 2e, f.

## Conclusion

This study conducted time-resolved imaging measurements of multi-ferroic BiFeO<sub>3</sub> and observed that antiferromagnetic magnon can propagate over hundreds of micrometers via magnon-photon coupling. The analyses based on Fourier transform provided wavenumber-resolved information that is inaccessible by conventional terahertz spectroscopy. Additionally, the refractive index of BiFeO<sub>3</sub>,  $n_{\text{THz}}$ , was extracted, and approximate value of the group velocity of magnon-polaritons was determined. Although the data were obtained using an unprocessed single crystal, future extensions such as coupling enhancement by a cavity opto-magnonic system are feasible<sup>37</sup>. Thus, our findings offer comprehensive insights into antiferromagnetic magnonics, polaritonics, and potential future applications.

## Methods

The 120  $\mu\text{m}$ -thick (111)<sub>pc</sub>-oriented BiFeO<sub>3</sub> single crystal was grown using the modified floating-zone method with laser diodes<sup>38</sup>, employing magnetic annealing in the  $[10\bar{1}]_{\text{pc}}$ -direction. Pump-probe imaging measurements for BiFeO<sub>3</sub> were conducted using femtosecond laser pulses. A Ti:sapphire regenerative amplifier with a repetition rate of 1 kHz and a pulse duration of  $\tau = 60 \text{ fs}$  was utilized to generate both the pump and probe pulses with a central wavelength of 800 nm. Additionally, the central wavelength of the pump pulse was converted to 1300 nm using an optical parametric amplifier. Linearly polarized pump pulses with a wavelength of 1300 nm were focused onto the sample as a line-shaped spot with a width of 20  $\mu\text{m}$  using a cylindrical lens. The pump pulses were directed to the left edge of the region depicted in Fig. 1b. The azimuthal angle of the pump polarization was adjusted using a half-wave plate. Additionally, the probe pulse with a wavelength of 800 nm was irradiated without focusing and with varying time delays relative to the pump pulse. The probe pulse was circularly polarized using a quarter-wave plate (QWP). Moreover, the probe pulses became elliptically polarized due to the sample's electro-optical or magneto-optical effect. The transmitted probe pulse was then transformed into an approximately linearly polarized pulse using another QWP. Additionally, the ellipticity change of the probe pulse was determined using a rotating-polarizer method with a wire grid analyzer and a CMOS camera<sup>39</sup>. An optical configuration sensitive to the off-diagonal component of the refractive index modulation was employed to selectively extract extraordinary electromagnetic waves<sup>33</sup>. The detection mechanism is explained in Supplementary Information III. The sample was placed in a cryostat for temperature-dependent measurements ranging from 5 to 300 K.

## Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author upon reasonable request.

Received: 25 February 2024; Accepted: 5 June 2024;

Published online: 02 July 2024

## References

- Kruglyak, V. V., Demokritov, S. O. & Grundler, D. Magnonics. *J. Phys. D: Appl. Phys.* **43**, 264001 (2010).
- Pirro, P., Vasyuchka, V. I., Serga, A. A. & Hillebrands, B. Advances in coherent magnonics. *Nat. Rev. Mater.* **6**, 1114–1135 (2021).
- Schneider, T. et al. Realization of spin-wave logic gates. *Appl. Phys. Lett.* **92**, 022505 (2008).
- Baltz, V. et al. Antiferromagnetic spintronics. *Rev. Mod. Phys.* **90**, 015005 (2018).
- Li, J. et al. Spin current from sub-terahertz-generated antiferromagnetic magnons. *Nature* **578**, 70–74 (2020).
- Chumak, A. V. et al. Advances in magnetics roadmap on spin-wave computing. *IEEE Trans. Magn.* **58**, 1–72 (2022).
- Hirsch, J. Spin Hall effect. *Phys. Rev. Lett.* **83**, 1834 (1999).
- Lebrun, R. et al. Tunable long-distance spin transport in a crystalline antiferromagnetic iron oxide. *Nature* **561**, 222–225 (2018).
- Wang, H. et al. Long-distance coherent propagation of high-velocity antiferromagnetic spin waves. *Phys. Rev. Lett.* **130**, 096701 (2023).
- Satoh, T. et al. Spin oscillations in antiferromagnetic NiO triggered by circularly polarized light. *Phys. Rev. Lett.* **105**, 077402 (2010).
- Němec, P., Fiebig, M., Kampfrath, T. & Kimel, A. V. Antiferromagnetic opto-spintronics. *Nat. Phys.* **14**, 229–241 (2018).
- Mills, D. L. & Burstein, E. Polaritons: the electromagnetic modes of media. *Rep. Prog. Phys.* **37**, 817 (1974).
- Grishunin, K. et al. Terahertz magnon-polaritons in TmFeO<sub>3</sub>. *ACS Photonics* **5**, 1375–1380 (2018).
- Boventer, I. et al. Antiferromagnetic cavity magnon polaritons in collinear and canted phases of hematite. *Phys. Rev. Appl.* **19**, 014071 (2023).
- Rovillain, P. et al. Electric-field control of spin waves at room temperature in multiferroic BiFeO<sub>3</sub>. *Nat. Mater.* **9**, 975–979 (2010).
- Fiebig, M., Lottermoser, T., Meier, D. & Trassin, M. The evolution of multiferroics. *Nat. Rev. Mater.* **1**, 1–14 (2016).
- Sivarajah, P. et al. THz-frequency magnon-phonon-polaritons in the collective strong-coupling regime. *J. Appl. Phys.* **125**, 213103 (2019).
- Hortensius, J. R. et al. Coherent spin-wave transport in an antiferromagnet. *Nat. Phys.* **17**, 1001–1006 (2021).
- Sosnowska, I., Neumaier, T. P. & Steichele, E. Spiral magnetic ordering in bismuth ferrite. *J. Phys. C: Solid State Phys.* **15**, 4835–4846 (1982).
- Burns, S. R., Paull, O., Juraszek, J., Nagarajan, V. & Sando, D. The experimentalist's guide to the cycloid, or noncollinear antiferromagnetism in epitaxial BiFeO<sub>3</sub>. *Adv. Mater.* **32**, 2003711 (2020).
- de Sousa, R. & Moore, J. E. Optical coupling to spin waves in the cycloidal multiferroic BiFeO<sub>3</sub>. *Phys. Rev. B* **77**, 012406 (2008).
- Cazayous, M. et al. Possible observation of cycloidal electromagnons in BiFeO<sub>3</sub>. *Phys. Rev. Lett.* **101**, 037601 (2008).
- Buhot, J. et al. Driving spin excitations by hydrostatic pressure in BiFeO<sub>3</sub>. *Phys. Rev. Lett.* **115**, 267204 (2015).
- Komandin, G. A. et al. Optical properties of BiFeO<sub>3</sub> ceramics in the frequency range 0.3–30.0 THz. *Phys. Solid State* **52**, 734–743 (2010).
- Talbayev, D. et al. Long-wavelength magnetic and magnetoelectric excitations in the ferroelectric antiferromagnet BiFeO<sub>3</sub>. *Phys. Rev. B* **83**, 094403 (2011).
- Skiadopoulou, S. et al. Spin and lattice excitations of a BiFeO<sub>3</sub> thin film and ceramics. *Phys. Rev. B* **91**, 174108 (2015).
- Matsubara, E., Mochizuki, T., Nagai, M. & Ashida, M. Self-polarized terahertz magnon absorption in a single crystal of BiFeO<sub>3</sub>. *Phys. Rev. B* **94**, 054426 (2016).
- Bialek, M., Magrez, A., Murk, A. & Ansermet, J.-P. Spin-wave resonances in bismuth orthoferrite at high temperatures. *Phys. Rev. B* **97**, 054410 (2018).
- Farkas, D. G. et al. Selection rules and dynamic magnetoelectric effect of the spin waves in multiferroic BiFeO<sub>3</sub>. *Phys. Rev. B* **104**, 174429 (2021).
- Khan, P., Kanamaru, M., Matsumoto, K., Ito, T. & Satoh, T. Ultrafast light-driven simultaneous excitation of coherent terahertz magnons and phonons in multiferroic BiFeO<sub>3</sub>. *Phys. Rev. B* **101**, 134413 (2020).
- Fishman, R. S. et al. Spin-induced polarizations and nonreciprocal directional dichroism of the room-temperature multiferroic BiFeO<sub>3</sub>. *Phys. Rev. B* **92**, 094422 (2015).
- Auston, D. H. Subpicosecond electro-optic shock waves. *Appl. Phys. Lett.* **43**, 713–715 (1983).
- Matsumoto, K. & Satoh, T. Selective imaging of the terahertz electric field of the phonon-polariton in LiNbO<sub>3</sub>. *Phys. Rev. B* **102**, 094313 (2020).
- Bialek, M., Ito, T., Rønnow, H. & Ansermet, J.-P. Terahertz-optical properties of a bismuth ferrite single crystal. *Phys. Rev. B* **99**, 064429 (2019).
- Yang, C., Wu, J., Xu, J., Nelson, K. A. & Werley, C. A. Experimental and theoretical analysis of THz-frequency, direction-dependent, phonon polariton modes in a subwavelength, anisotropic slab waveguide. *Opt. Express* **18**, 26351–26364 (2010).
- Mahmoud, A. et al. Introduction to spin wave computing. *J. Appl. Phys.* **128**, 161101 (2020).
- Parvini, T. S., Bittencourt, V. A. S. V. & Kusminskiy, S. V. Antiferromagnetic cavity optomagnonics. *Phys. Rev. Res.* **2**, 022027 (2020).
- Ito, T., Ushiyama, T., Yanagisawa, Y., Kumai, R. & Tomioka, Y. Growth of highly insulating bulk single crystals of multiferroic BiFeO<sub>3</sub> and their inherent internal strains in the domain-switching process. *Cryst. Growth Des.* **11**, 5139–5143 (2011).
- Yoshimine, I. et al. Phase-controllable spin wave generation in iron garnet by linearly polarized light pulses. *J. Appl. Phys.* **116**, 043907 (2014).

## Acknowledgements

We thank K. T. Yamada for the valuable discussions and technical support.

## Author contributions

T.S. conceived and supervised the study. R.K., K.M., and T.S. performed the experiments and analyzed the data. T.I. fabricated samples. R.K. and T.S. wrote the manuscript. All authors discussed the results and provided comments on the manuscript.

## Competing interests

The authors declare no competing interests.

## Additional information

**Supplementary information** The online version contains supplementary material available at <https://doi.org/10.1038/s44306-024-00034-3>.

**Correspondence** and requests for materials should be addressed to Takuya Satoh.

**Reprints and permissions information** is available at <http://www.nature.com/reprints>

**Publisher's note** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

**Open Access** This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit <http://creativecommons.org/licenses/by/4.0/>.

© The Author(s) 2024